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# Numerical solution of nonlinear Fredholm-Hammerstein integral equations with logarithmic kernel by spline quasi-interpolating projectors

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# Abstract

Nonlinear Fredholm-Hammerstein integral equations with logarithmic kernel are here taken into account and numerically solved by spline quasi-interpolating projectors based collocation and Kulkarni methods, both in their basic and iterated versions. Theoretical analysis of discretization error and convergence order is provided, together with numerical results validating the estimates obtained under the hypothesis of sufficiently smooth solutions. Finally, some results in case of less regular solutions show the robustness of the proposed approach even in a non smooth framework.

Keywords: Fredholm-Hammerstein integral equation; Logarithmic kernel; Spline quasi-interpolating projector; Kulkarni

#### 1. Introduction

Many problems in the applied sciences lead to mathematical models described by nonlinear integral equations. For instance, the Fredholm-Hammerstein integral equations appear in a variety of applications in many fields including continuum mechanics, potential theory, geophysics, electromagnetic fluid dynamics, antenna synthesis problem, communication theory, mathematical economics, population genetics, radiation, the particle transport problems of astrophysics and reactor theory, fluid mechanics (see e.g. [13, 22] and references therein). In the nineties, many papers appeared and handled these integral equations. Among them, significant results have been proved in [13, 25, 26, 27, 32] and, in the literature, it is highlighted that the convergence orders of numerical methods are affected by possible singular behaviour of the solutions near the domain boundary.

On the other side, various practical engineering problems, such as for instance the analysis of linear transport equation in slab geometry [31] or the study of steady potential flow past obstacles [29], such as around airfoils by Boundary Element Methods (BEMs) [18], give rise to integral equations with weakly singular kernels of logarithmic type.

Here, the focus is on the numerical solution of nonlinear Fredholm-Hammerstein integral equations with logarithmic kernel, appearing for instance in the integral reformulation of two-dimensional elliptic boundary value problems with a nonlinear boundary condition. For example, in [35], several applications of integral equations are presented and for the considered model the exterior boundary value problem for the two-dimensional Helmholtz equation is proposed. In general, it is difficult to find the exact solution of these type of integral equations and hence to obtain approximate solutions is often mandatory. The problem has recently received increasing attention and several papers in the literature investigate the topic, see e.g. [10, 11, 24, 30].

Classical methods to search an approximate solution of the above kind of problems are the projection ones, such as the Galerkin and collocation; recently, a modified projection method, based on a sequence of orthogonal or interpolatory projectors onto finite dimensional subspaces, usually spaces of piecewise polynomials

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of degree d at most continuous, has been proposed by Kulkarni [24, 28].

On the other side, the use of the spline quasi-interpolation has been proved to work well for the approximation of the solution of Fredholm integral equations, also in the bivariate case (see e.g. [5, 6, 7, 8, 16, 19, 20, 21]). Therefore, here the novelty w.r.t. the state of the art lies in the use of spline Quasi-Interpolating Projectors (QIPs) in the space of splines of degree d and smoothness  $C^{d-1}$ , as approximation tools in the framework of collocation and Kulkarni discretization techniques for the solution of nonlinear Fredholm-Hammerstein integral equations with logarithmic kernel, under the hypothesis that the data are assigned such that the problems admit sufficiently regular solutions (for the interested reader, some examples of integral equations with logarithmic kernel and regular solutions can be found in [12, 4, 35, 9, 14]).

This in order to theoretically study the convergence order of the proposed numerical approaches. In fact, the regularity of the solution is needed to take advantage of the use of the spline QIPs, since their convergence properties, on which the convergence order of the employed methods relies, are based on the regularity of the function to be approximated, i.e. in this framework, the solution of the integral equation [17].

We underline that this work can be conceived as a prosecution of [21, 19], where spline QIPs were introduced for the discretization of linear and nonlinear Fredholm integral equations, in both cases with smoother kernels. We also remark that the use of smooth splines allows an advantage from the computational point of view with respect to the growth of the dimension of linear systems to be solved in the construction of the approximate solutions, as explained at the end of Section 4.1.

The paper is structured as follows: in the next Section 2 the problem at hand is described, while in Section 3 spline QIPs are briefly recalled. Then in Section 4, spline QIP based collocation and Kulkarni numerical methods, in basic and iterated versions, are illustrated and their discretization error, with subsequent convergence order, is studied. Section 5 contains numerical results that validate the theoretical estimates. Moreover, the reader will find some results in case of less regular solutions, which show the robustness of the proposed approach even in a non smooth framework. Conclusions are summarized, together with future research lines, in Section 6. For readers' convenience, basic quadrature rules for logarithmic kernel are resumed from [2] in Appendix A, together with the description of their iterated use as implemented for the numerical simulations, while an extension of a theoretical result in [19], needed to prove some error estimates, is given in Appendix B.

#### 2. Fredholm-Hammerstein integral equations with logarithmic kernel

In this paper we consider Fredholm-Hammerstein integral equations of the second kind of the form

$$x - K(x) = f, (2.1)$$

where K is the integral operator

$$K(x)(s) := \int_0^1 k(s,t)\psi(t,x(t))dt, \qquad s \in \mathcal{I} := [0,1], \qquad x \in X := C(\mathcal{I}),$$

 $k(s,t) = \log |s-t|$  is the logarithmic kernel taken into account, f and  $\psi$  are known functions and x is the unknown solution to be determined.

We assume throughout this paper the following conditions on f and  $\psi$ :

- 1.  $f \in X$ ,
- 2. for  $x \in X$ ,  $\psi(\cdot, x(\cdot)) \in X$ ,
- 3.  $\psi(t,x(t))$  is bounded and continuous over  $\mathcal{I} \times \mathbb{R}$  and Lipschitz continuous in x, i.e. there exists a constant  $c_1 > 0$  for which  $|\psi(t, u) - \psi(t, v)| \le c_1 |u - v|$ ,  $\forall u, v \in \mathbb{R}$ ,
- 4. the partial derivative  $\frac{\partial \psi}{\partial x}(t,x(t))$  of  $\psi$  with respect to the second variable exists and we assume it is Lipschitz continuous with respect to the second variable, i.e. there exists  $c_2 > 0$  such that

$$\left| \frac{\partial \psi}{\partial x}(t, u) - \frac{\partial \psi}{\partial x}(t, v) \right| \le c_2 |u - v|, \quad \forall u, v \in \mathbb{R},$$

5. 
$$B = \sup_{t \in \mathcal{I}} \left| \frac{\partial \psi}{\partial x}(t, x(t)) \right| < \infty.$$

The Fréchet derivative of K is given by [30]

$$(K'(x)q)(s) = \int_0^1 k(s,t) \frac{\partial \psi}{\partial x}(t,x(t))q(t)dt$$
 (2.2)

(for a general definition see [15]).

Let us note that, for any  $s \in \mathcal{I}$ , we have

$$\int_{0}^{1} |k(s,t)| dt \le p_{1} < \infty, \quad \text{where } p_{1} = 1.7$$
 (2.3)

and if  $c_2 p_1 < 1$  equation (2.1) admits a unique solution [25], denoted in the following by  $\varphi$ . Furthermore,

$$\lim_{t \to \tau} ||k(\cdot, t) - k(\cdot, \tau)||_1 = 0, \quad \tau \in \mathcal{I}.$$

Finally, we define the operator T by T(u) = f + K(u), so that (2.1) can be written as x = T(x).

# 3. Spline quasi-interpolating projectors

In order to make the paper self-contained, in this section we recall definition and properties of the spline QIPs proposed in [19].

Let us consider the space of splines of degree d and class  $C^{d-1}(\mathcal{I})$  on the uniform knot partition  $\mathcal{T}_n := \{t_i = ih, 0 \le i \le n\}$ , with h = 1/n. We denote such a space by  $X_n$  and here we consider QIPs on  $X_n$  of the following form

$$\pi_n x := \sum_{i=1}^N \lambda_i(x) B_i, \qquad x \in X$$
(3.1)

where:

- $N := \dim(X_n) = n + d;$
- $\mathcal{T}_n^e := \mathcal{T}_n \cup \{t_{-d} = \ldots = t_0 = 0; 1 = t_n = \ldots = t_{n+d}\}$  is the usual extended knot sequence associated with  $\mathcal{T}_n$ ;
- $\{B_i\}_{i=1}^N$  are the B-splines with support supp $B_i = [t_{i-d-1}, t_i]$  on  $\mathcal{T}_n^e$ , forming a basis for  $X_n$ ;
- $\left\{\lambda_{i}\right\}_{i=1}^{N}$  are point coefficient functionals of the form

$$\lambda_i(x) := \sum_{j=2(i-d-1)}^{2i} \sigma_{i,j} x_j, \qquad x_j := x(\xi_j)$$
 (3.2)

based on the quasi-interpolation nodes  $\{\xi_j\}_{j=0}^{2n}$  (the QI nodes involved in (3.2) are inside supp  $B_i$ ) with

$$\left\{ \begin{array}{ll} \xi_{2i} := t_i, & 0 \leq i \leq n \\ \xi_{2i-1} := s_i := \frac{1}{2}(t_{i-1} + t_i) & 1 \leq i \leq n \end{array} \right.$$

and the  $\sigma_{i,j}$ 's chosen such that  $\pi_n x = x$ , for all  $x \in X_n$ .

The QIP  $\pi_n$  can also be written in the quasi-Lagrange form, by means of the so-called fundamental functions, given by linear combination of B-splines, according to (3.2)

$$\pi_n x = \sum_{i=0}^{2n} x_i L_i. \tag{3.3}$$

Here we recall some properties of the QIP  $\pi_n$ :

•  $\pi_n$  is bounded, i.e.

$$\|\pi_n\|_{\infty} := \sup_{x \in X, x \neq 0} \frac{\|\pi_n x\|_{\infty}}{\|x\|_{\infty}} < \infty$$

because the  $\lambda_i$  are continuous linear functionals;

• from classical results in approximation theory

$$||x - \pi_n x||_{\infty} \le C \operatorname{dist}(x, X_n), \qquad C := 1 + ||\pi_n||_{\infty};$$

• for  $x \in C^j(\mathcal{I})$ , there exists a constant  $\overline{C}_j$ , depending on C and j, such that,

$$||x - \pi_n x||_{\infty} \le \overline{C}_j h^j \omega(x^{(j)}, h), \text{ with } 0 \le j \le d,$$

where  $\omega$  is the modulus of continuity of  $x^{(j)}$ . Moreover, if  $x \in C^{d+1}(\mathcal{I})$  we have

$$||x - \pi_n x||_{\infty} = O(h^{d+1}).$$
 (3.4)

These results are deduced from Jackson type theorem for splines [17, 34].

Examples of spline QIPs of the form (3.1) can be found in [19]. They are denoted by  $Q_2$  and  $Q_3$  and are used in the numerical tests here proposed. The operators  $Q_2$  is defined in  $\mathcal{S}_2^1(\mathcal{I}, \mathcal{T}_n)$  and  $Q_3$  is defined on  $\mathcal{S}_3^2(\mathcal{I}, \mathcal{T}_n)$ . For both of them, the following inequality holds:

$$\sup_{p} ||Q_j|| \le p_j, \quad j = 2, 3,$$

where  $p_i$  are suitable positive real constants [21].

The following theorems present some interesting properties of the projectors  $\pi_n$ , in case of even degree d (for details see [19]).

**Theorem 1.** Let  $\pi_n$  be a QIP on  $X_n$  of kind (3.1) and let the degree d be even. If the functionals  $\lambda_i$ ,  $i = d+1, \ldots, n$ , are such that the values  $\sigma_{i,j}$  in (3.2), associated to QI nodes symmetric w.r.t. the center of  $supp B_i$ , are equal, then

$$\int_{t_{i-1}}^{t_i} (\pi_n m_{d+1}(t) - m_{d+1}(t)) dt = 0, \qquad i = d+1, \dots, n-d, \qquad m_{d+1}(t) = t^{d+1}.$$

It is interesting to consider QIPs for which Theorem 1 is valid also in the case of odd degree, as it is the case for the QIP  $Q_3$  (see Appendix B).

**Theorem 2.** If Theorem 1 holds, for any function  $g \in W^{1,1}$  (i.e. with  $||g'||_1$  bounded) and any function x such that  $||x^{(d+2)}||_{\infty}$  is bounded, there results

$$\left| \int_0^1 g(t)(\pi_n x(t) - x(t)) dt \right| = O(h^{d+2}).$$

For the considered QIPs  $Q_2$  and  $Q_3$ , Theorem 2 holds.

### 4. Spline projection methods

Given a spline QIP operator  $\pi_n: X \to X_n$ , defined as in Section 3, we introduce in the following two projection methods based on it. The first one is considered for its good numerical performance, while the second one, which is standard, is taken into account for comparison purposes:

1. QIP spline Kulkarni's type method. K is approximated by

$$K_n^k := \pi_n K + K \pi_n - \pi_n K \pi_n \tag{4.1}$$

in (2.1) and the approximate equation is

$$\varphi_n^k - K_n^k(\varphi_n^k) = f. \tag{4.2}$$

Defining  $T_n^k$  by  $T_n^k(u) = f + K_n^k(u)$ , then (4.2) can be written as  $\varphi_n^k = T_n^k(\varphi_n^k)$ . We also consider the iterated version of this Kulkarni's type method. By using the approximation  $\varphi_n^k$ and the equation (2.1) we get

$$\tilde{\varphi}_n^k = K(\varphi_n^k) + f,$$

where  $\tilde{\varphi}_n^k$  is the approximation of the solution by this iterated method.

2. QIP spline collocation method. K is approximated by  $K_n^c := \pi_n K \pi_n$  and the right hand side f by  $\pi_n f$ in (2.1), obtaining the approximate equation

$$\varphi_n^c - \pi_n K(\varphi_n^c) = \pi_n f. \tag{4.3}$$

Also in this case we consider the iterated version of the method. By using  $\varphi_n^c$  and (2.1) we get

$$\tilde{\varphi}_n^c = K(\varphi_n^c) + f,$$

where  $\tilde{\varphi}_n^c$  is the approximation of the solution by this iterated method.

4.1. Construction of the approximate solutions

Starting from equations (4.2) and (4.3), in this section we construct the corresponding approximate solutions.

1. QIP spline Kulkarni's type method.

We consider definition (4.1) and equation (4.2), we project them in  $X_n$  by using the spline QIP  $\pi_n$ , and we join them, obtaining

$$\pi_n \varphi_n^k - \pi_n K(\varphi_n^k) = \pi_n f. \tag{4.4}$$

Combining (4.4) with (4.2) we obtain

$$\varphi_n^k = K(\pi_n \varphi_n^k) - \pi_n K(\pi_n \varphi_n^k) + \pi_n \varphi_n^k - \pi_n f + f. \tag{4.5}$$

Now we define  $\psi_n := \pi_n \varphi_n^k$ ; from (4.4) we have

$$\psi_n - \pi_n K(\varphi_n^k) = \pi_n f \tag{4.6}$$

and from (4.5) we obtain

$$\varphi_n^k = \psi_n + (I - \pi_n)(K(\psi_n) + f). \tag{4.7}$$

Replacing (4.7) in (4.6) we finally have

$$\psi_n - \pi_n K(\psi_n + (I - \pi_n)(K(\psi_n) + f)) = \pi_n f, \tag{4.8}$$

where the unknown  $\psi_n$  by its definition lies in  $X_n$ .

In order to find  $\psi_n$ , we define the functional

$$F_n(y) = y - \pi_n K(y + (I - \pi_n)(K(y) + f)) - \pi_n f, \qquad y \in X_n, \tag{4.9}$$

with Fréchet derivative

$$F'_n(y)q = q - \pi_n K'(y + (I - \pi_n)(K(y) + f))(I + (I - \pi_n)K'(y))q. \tag{4.10}$$

We notice that (4.8) is equivalent to

$$F_n(\psi_n) = 0,$$

which is iteratively solved by Newton-Kantorovich method, an extension to functional spaces of the classical Newton method for the numerical solutions of nonlinear equations in one variable (see e.g. [33] for details).

Let  $\psi_n^{(0)}$  be the initial approximation needed by the method. The iterates  $\psi_n^{(r)}$ ,  $r=0,1,2,\ldots$ , are given by

$$\psi_n^{(r+1)} = \psi_n^{(r)} - [F_n'(\psi_n^{(r)})]^{-1} F_n(\psi_n^{(r)})$$

or, equivalently

$$F_n'(\psi_n^{(r)})\psi_n^{(r+1)} = F_n'(\psi_n^{(r)})\psi_n^{(r)} - F_n(\psi_n^{(r)}). \tag{4.11}$$

By using (4.9) and (4.10), and also (4.7), the equation (4.11) can be written in this way

$$\psi_n^{(r+1)} - \pi_n K'(\varphi_n^{(r)}) \psi_n^{(r+1)} - \pi_n K'(\varphi_n^{(r)}) (I - \pi_n) K'(\psi_n^{(r)}) \psi_n^{(r+1)}$$

$$= \pi_n (K(\varphi_n^{(r)}) + f) - \pi_n K'(\varphi_n^{(r)}) \psi_n^{(r)} - \pi_n K'(\varphi_n^{(r)}) (I - \pi_n) K'(\psi_n^{(r)}) \psi_n^{(r)}. \tag{4.12}$$

Recalling that  $\psi_n^{(r)} \in X_n$ , we can express it as a linear combination of B-splines

$$\psi_n^{(r)} = \sum_{j=1}^N x_n^{(r)}(j) B_j \qquad x_n^{(r)} \in \mathbb{R}^N.$$
(4.13)

After some algebra, we can write (4.12) in this way:

$$x_n^{(r+1)}(i) - \sum_{j=1}^N x_n^{(r+1)}(j)\lambda_i(K'(\varphi_n^{(r)})B_j) - \sum_{j=1}^N x_n^{(r+1)}(j)\lambda_i(K'(\varphi_n^{(r)})(I - \pi_n)K'(\psi_n^{(r)})B_j)$$

$$= \lambda_i(K(\varphi_n^{(r)})) + \lambda_i(f) - \sum_{j=1}^N x_n^{(r)}(j)\lambda_i(K'(\varphi_n^{(r)})B_j) - \sum_{j=1}^N x_n^{(r)}(j)\lambda_i(K'(\varphi_n^{(r)})(I - \pi_n)K'(\psi_n^{(r)})B_j),$$

 $i=1,\ldots,N$ . This is a linear system of size N, whose matrix form is

$$(I - \Lambda_n^{(r)} - \Xi_n^{(r)}) x_n^{(r+1)} = \delta_n^{(r+1)}, \tag{4.14}$$

where, for  $i, j = 1, \dots, N$ 

$$\Lambda_n^{(r)}(i,j) := \lambda_i (K'(\varphi_n^{(r)}) B_j)$$

$$\Xi_n^{(r)}(i,j) := \lambda_i (K'(\varphi_n^{(r)}) (I - \pi_n) K'(\psi_n^{(r)}) B_j)$$

$$\delta_n^{(r)}(i) := \lambda_i (K(\varphi_n^{(r)})) + \lambda_i (f) - (\Lambda_n^{(r)} x_n^{(r)}) (i) - (\Xi_n^{(r)} x_n^{(r)}) (i).$$

By solving the system (4.14), due to the non singularity of the related matrix, we get the vector  $x_n^{(r+1)}$ . Using this we can calculate  $\psi_n^{(r+1)}$  from (4.13). The approximate solution at the (r+1) iteration is  $\varphi_n^{(r+1)}$ , which can be constructed using (4.7).

# 2. QIP spline collocation method.

We start considering equation (4.3). We recall that in this method  $\varphi_n^c \in X_n$ . Since the equation is nonlinear, we solve it by Newton-Kantorovich method. In this case the functional needed in order to set the method is given by

$$F_n(y) = y - \pi_n K(y) - \pi_n f$$
  $y \in X_n$ 

and its Fréchet derivative by

$$F'_n(y)q = q - \pi_n K'(y)q.$$

Consequently, we notice that the equation

$$F_n(\varphi_n^c) = 0$$

is equivalent to (4.3), so the iteration of the Newton-Kantorovich method (4.11) is given by

$$\varphi_n^{(r+1)} - \pi_n K'(\varphi_n^{(r)}) \varphi_n^{(r+1)} = \pi_n (K(\varphi_n^{(r)}) + f) - \pi_n K'(\varphi_n^{(r)}) \varphi_n^{(r)}. \tag{4.15}$$

Since  $\varphi_n^{(r)} \in X_n$ , we can write

$$\varphi_n^{(r)} = \sum_{j=1}^N x_n^{(r)}(j) B_j, \qquad x_n^{(r)} \in \mathbb{R}^N.$$
(4.16)

From (4.15), after some algebra, we obtain

$$x_n^{(r+1)}(i) - \sum_{j=1}^N x_n^{(r+1)}(j) \lambda_i \left[ K'(\varphi_n^{(r)}) B_j \right] = \lambda_i \left[ K(\varphi_n^{(r)}) \right] + \lambda_i(f) - \sum_{j=1}^N x_n^{(r)}(j) \lambda_i \left[ K'(\varphi_n^{(r)}) B_j \right],$$

 $i=1,\ldots,N$ , that is a linear system of size N, whose matrix form is given by

$$(I - \Phi_n^{(r)}) x_n^{(r+1)} = \omega_n^{(r)}, \tag{4.17}$$

where, for  $i, j = 1, \dots, N$ 

$$\Phi_n^{(r)}(i,j) = \lambda_i \left( K'(\varphi_n^{(r)}) B_j \right)$$

$$\omega_n^{(r)}(i) = \lambda_i \big(K(\varphi_n^{(r)})\big) + \lambda_i(f) - \big(\Phi_n^{(r)} x_n^{(r)}\big)(i).$$

By solving the system (4.17), due to the non singularity of the related matrix, we get the vector  $x_n^{(r+1)}$ . Using this we can calculate  $\varphi_n^{(r+1)}$  from (4.16), which is the approximate solution at the (r+1) iteration.

We remark that classical methods for the solution of the above kind of problems are the projection ones based on a sequence of orthogonal or interpolatory projectors, usually onto spaces of piecewise polynomials of degree d at most continuous. In this case the dimension of the linear systems is related to the product between the number of subintervals n and the degree d. Instead, in our approach the dimension of the linear systems is related to the sum between the number of subintervals n and the degree d and therefore we have an advantage from the computational point of view for increasing values of n.

# 4.2. Convergence of the methods

In this section we study the convergence of the spline projection methods (4.2) and (4.3) and their iterated version

Concerning the existence and uniqueness of the approximate solutions  $\varphi_n^k$  and  $\varphi_n^c$ , we can refer to the general results given in [30] and [26], respectively, that also hold for the considered spline QIPs.

First of all we prove the following lemma.

**Lemma 3.** Let  $\varphi \in C^{d+2}(\mathcal{I})$  be an isolated solution of (2.1) and assume that 1 is not an eigenvalue of  $K'(\varphi)$  and let  $\pi_n : X \to X_n$  be a spline QIP of kind (3.1) for which Theorem 1 is valid. Then

$$||K(\pi_n \varphi) - K(\varphi)||_{\infty} = O(h^{d+2} \log(h)).$$

**Proof.** Consider

$$|(K(\pi_n\varphi) - K(\varphi))(s)| = \left| \int_0^1 \log|s - t| \left[ \psi(t, \pi_n\varphi(t)) - \psi(t, \varphi(t)) \right] dt \right|$$

$$\leq \left| \int_0^1 \log|s - t| \left[ \frac{\partial \psi}{\partial \varphi}(t, \varphi(t) + \theta_1(\pi_n\varphi - \varphi)(t)) - \frac{\partial \psi}{\partial \varphi}(t, \varphi(t)) \right] (\pi_n\varphi - \varphi)(t) dt \right|$$

$$+ \left| \int_0^1 \log|s - t| \frac{\partial \psi}{\partial \varphi}(t, \varphi(t)) (\pi_n\varphi - \varphi)(t) dt \right| = |\clubsuit| + |\spadesuit|,$$

with  $0 < \theta_1 < 1$ . Then, since  $\frac{\partial \psi}{\partial x}(t, x(t))$  is Lipschitz continuous with constant  $c_2$ , from (2.3) and (3.4) we have

$$| \clubsuit | \leq \int_0^1 |\log|s - t| \left| \frac{\partial \psi}{\partial \varphi}(t, \varphi(t) + \theta_1(\pi_n \varphi - \varphi)(t)) - \frac{\partial \psi}{\partial \varphi}(t, \varphi(t)) \right| \left| (\pi_n \varphi - \varphi)(t) \right| dt$$

$$\leq c_2 \theta_1 \int_0^1 |\log|s - t| \left| (\pi_n \varphi - \varphi)(t) \right|^2 dt \leq c_2 \theta_1 p_1 \left\| \pi_n \varphi - \varphi \right\|_{\infty}^2 = O(h^{2d+2}).$$

Now we consider the second term  $| \spadesuit |$ . Defining  $g_s(t) := \log |s-t| \frac{\partial \psi}{\partial \varphi}(t, \varphi(t))$  and choosing  $\phi_s(t)$  the polynomial of degree less than or equal to n as in Lemma 3.7 in [30] (see also [34, p. 92]), such that  $||g_s - \phi_s||_1 = O(h \log(h))$ , from Theorem 2 and again (3.4), we get

$$| \spadesuit | = \left| \int_0^1 g_s(t) (\pi_n \varphi - \varphi)(t) dt \right| \le \left| \int_0^1 (g_s - \phi_s)(t) (\pi_n \varphi - \varphi)(t) dt \right| + \left| \int_0^1 \phi_s(t) (\pi_n \varphi - \varphi)(t) dt \right|$$

$$\le \|g_s - \phi_s\|_1 \|\pi_n \varphi - \varphi\|_{\infty} + O(h^{d+2}) = O(h^{d+2} \log(h)) + O(h^{d+2}) = O(h^{d+2} \log(h))$$

and the thesis follows.  $\Box$ 

Now, we consider the QIP spline collocation method and its iterated version and we prove the following results.

**Theorem 4.** Let  $\varphi \in C^{d+1}(\mathcal{I})$  be an isolated solution of (2.1) and assume that 1 is not an eigenvalue of  $K'(\varphi)$ . Let  $\pi_n : X \to X_n$  be a spline QIP of kind (3.1). Then (4.3) has a unique solution  $\varphi_n^c \in B(\varphi, \delta) = \{x : \|x - \varphi\|_{\infty} < \delta\}$  for some  $\delta > 0$  and for sufficiently large n. Moreover

$$\|\varphi_n^c - \varphi\|_{\infty} = O(h^{d+1}).$$

**Proof.** We follow the lines of the proof of Theorem 2.1 in [25], reaching the inequality

$$\|\varphi_n^c - \varphi\|_{\infty} \le c_3 \|\pi_n \varphi - \varphi\|_{\infty},$$

for a suitable constant  $c_3$ . Using (3.4) the thesis holds.

**Theorem 5.** Let  $\varphi \in C^{d+2}(\mathcal{I})$  be an isolated solution of (2.1) and assume that 1 is not an eigenvalue of  $K'(\varphi)$ . Let  $\pi_n : X \to X_n$  be a spline QIP of kind (3.1) for which Theorem 1 is valid. Let  $\tilde{\varphi}_n^c$  be the iterated approximation of the spline collocation method. Then, there holds

$$\|\tilde{\varphi}_n^c - \varphi\|_{\infty} = O(h^{d+2}\log(h)).$$

**Proof.** Following the same path of reasoning used in the proof of Lemma 3, from the definition of  $\tilde{\varphi}_n^c$  we have

$$\tilde{\varphi}_{n}^{c} - \varphi = K(\varphi_{n}^{c}) - K(\varphi) = \left[K'\left(\varphi + \theta_{2}\left(\varphi_{n}^{c} - \varphi\right)\right) - K'\left(\varphi\right)\right]\left(\varphi_{n}^{c} - \varphi\right) + K'\left(\varphi\right)\left(\varphi_{n}^{c} - \varphi\right),\tag{4.18}$$

with  $0 < \theta_2 < 1$ . Taking into account that  $\pi_n \tilde{\varphi}_n^c = \varphi_n^c$ , we have

$$\tilde{\varphi}_{n}^{c} - \varphi = \left[K'\left(\varphi + \theta_{2}\left(\varphi_{n}^{c} - \varphi\right)\right) - K'\left(\varphi\right)\right]\left(\varphi_{n}^{c} - \varphi\right) + K'\left(\varphi\right)\left(\pi_{n}\left(\tilde{\varphi}_{n}^{c} - \varphi\right)\right) + K'\left(\varphi\right)\left(\pi_{n}\varphi - \varphi\right)\right)$$

and consequently

$$[I - K'(\varphi)\pi_n](\tilde{\varphi}_n^c - \varphi) = [K'(\varphi + \theta_2(\varphi_n^c - \varphi)) - K'(\varphi)](\varphi_n^c - \varphi) + K'(\varphi)(\pi_n\varphi - \varphi).$$

Following the theory of the collectively compact operators [26, 27], which can be applied thanks to the hypothesis reported in Section 2, we have that the inverse of the operator  $[I - K'(\varphi) \pi_n]$  exists with bounded infinity norm. Therefore

$$\|\tilde{\varphi}_{n}^{c} - \varphi\|_{\infty} \leq c_{4} \|\varphi_{n}^{c} - \varphi\|_{\infty}^{2} + c_{5} \|K'(\varphi)(\pi_{n}\varphi - \varphi)\|_{\infty},$$

for suitable constants  $c_4$ ,  $c_5$ . Following the proof of Lemma 3 and using the results there obtained, there exists a constant  $c_6$  such that

$$\|K'(\varphi)(\pi_n\varphi - \varphi)\|_{\infty} \le c_6 h \log(h) \|\pi_n\varphi - \varphi\|_{\infty} + \sup_{s \in \mathcal{I}} \left| \int_0^1 \phi_s(t)(\pi_n\varphi - \varphi)(t) dt \right|$$

where  $\phi_s(t)$  is the polynomial of degree less than or equal to n already chosen in Lemma 3. Hence, from (3.4) and Theorem 2, the thesis follows.

**Remark.** The rigorous proof of Theorem 5 can be substituted by an alternative one, based on a combination of theoretical steps and an assumption supported by an intensive numerical testing, which can be outlined as follows. Treating the first term in the right-hand side of (4.18) similarly as  $\clubsuit$  and the second term similarly as  $\spadesuit$ , we obtain

$$\|\tilde{\varphi}_n^c - \varphi\|_{\infty} \le c_7 \|\varphi_n^c - \varphi\|_{\infty}^2 + c_8 h \log(h) \|\varphi_n^c - \varphi\|_{\infty} + \sup_{s \in \mathcal{I}} \left| \int_0^1 \phi_s(t) (\varphi_n^c - \varphi)(t) dt \right|, \tag{4.19}$$

for suitable constants  $c_7$ ,  $c_8$  and  $\phi_s(t)$  the polynomial of degree less than or equal to n already chosen in Lemma 3. On the basis of several numerical evidences, we can conjecture that the last term of (4.19) decays as  $O(h^{d+2})$ , i.e. with an extra order with respect to the error  $\|\varphi_n^c - \varphi\|_{\infty}$ . Therefore, from this fact and from Theorem 4, (4.19) gives the thesis.

Considering the QIP spline Kulkarni method, we obtain the following result.

**Theorem 6.** Let  $\varphi \in C^{d+2}(\mathcal{I})$  be an isolated solution of (2.1) and assume that 1 is not an eigenvalue of  $K'(\varphi)$ . Let  $\pi_n : X \to X_n$  be a spline QIP of kind (3.1) for which Theorem 1 is valid. Then (4.2) has a unique solution  $\varphi_n^k \in B(\varphi, \delta) = \{x : ||x - \varphi||_{\infty} < \delta\}$  for some  $\delta > 0$  and for sufficiently large n. Moreover, there exists a constant 0 < q < 1, independent of n, such that

$$\frac{\alpha_n}{1+q} \le \left\| \varphi_n^k - \varphi \right\|_{\infty} \le \frac{\alpha_n}{1-q} \tag{4.20}$$

where  $\alpha_n := \|[I - (T_n^k)'(\varphi)]^{-1}(T_n^k(\varphi) - T(\varphi))\|_{\infty}$ . Further

$$\left\|\varphi_n^k - \varphi\right\|_{\infty} = O(h^{d+2}\log(h)). \tag{4.21}$$

**Proof.** Following the theory of the collectively compact operators [26, 27], which can be applied thanks to the hypothesis reported in Section 2, we have that  $(I - (T_n^k)'(\varphi))$  is invertible and

$$\|[I - (T_n^k)'(\varphi)]^{-1}\|_{\infty} \le c_9.$$

Using this inequality and Lemma 3.4 of [30], after some algebra, we can say that hypothesis of Theorem 3.2 of [30] hold, so (4.20) is proved.

From the last inequality of (4.20) we can write

$$\|\varphi_{n}^{k} - \varphi\|_{\infty} \leq \frac{\alpha_{n}}{1 - q} \leq \frac{c_{9}}{1 - q} \|T_{n}^{k}(\varphi) - T(\varphi)\|_{\infty} \leq \frac{c_{9}}{1 - q} \|(I - \pi_{n})(K(\pi_{n}\varphi) - K(\varphi))\|_{\infty}$$
$$\leq c_{10} \|K(\pi_{n}\varphi) - K(\varphi)\|_{\infty},$$

for a suitable constant  $c_{10}$ . Applying Lemma 3, (4.21) holds.

Finally, for the iterated version of the QIP spline Kulkarni method, we give the following result, whose claim has been proved only by a combination of rigorous theoretical steps and an assumption supported by an intensive numerical testing, following the path of reasoning as in the Remark written at the end of Theorem 5.

**Proposition 7.** Let  $\varphi \in C^{d+2}(\mathcal{I})$  be an isolated solution of (2.1) and assume that 1 is not an eigenvalue of  $K'(\varphi)$ . Let  $\pi_n : X \to X_n$  be a spline QIP of kind (3.1) for which Theorem 1 is valid. Let  $\tilde{\varphi}_n^k$  be the iterated approximation of the Kulkarni's type method. Then, there holds

$$\|\tilde{\varphi}_n^k - \varphi\|_{\infty} = O(h^{d+3}(\log(h))^2). \tag{4.22}$$

**Proof.** From the definition of  $\tilde{\varphi}_n^k$ , we have

$$\tilde{\varphi}_{n}^{k}-\varphi=K(\varphi_{n}^{k})-K(\varphi)=\left[K'\left(\varphi+\theta_{3}\left(\varphi_{n}^{k}-\varphi\right)\right)-K'\left(\varphi\right)\right]\left(\varphi_{n}^{k}-\varphi\right)+K'\left(\varphi\right)\left(\varphi_{n}^{k}-\varphi\right),$$

with  $0 < \theta_3 < 1$ . Following the same path of reasoning used in the proof of Lemma 3, we obtain

$$\left\|\tilde{\varphi}_{n}^{k} - \varphi\right\|_{\infty} \leq c_{11} \left\|\varphi_{n}^{k} - \varphi\right\|_{\infty}^{2} + c_{12} h \log(h) \left\|\varphi_{n}^{k} - \varphi\right\|_{\infty} + \sup_{s \in \mathcal{I}} \left|\int_{0}^{1} \phi_{s}(t)(\varphi_{n}^{k} - \varphi)(t)dt\right|, \tag{4.23}$$

for suitable constants  $c_{11}$ ,  $c_{12}$  and  $\phi_s(t)$  the polynomial of degree less than or equal to n chosen in Lemma 3. On the basis of several numerical evidences, we can conjecture that the last term of (4.23) decays as  $O(h^{d+3}\log(h))$ , i.e. with an extra order with respect to the error  $\|\varphi_n^k - \varphi\|_{\infty}$ . Therefore, from this fact and from Theorem 6, (4.23) gives the thesis.

From Theorem 6 and Proposition 7, we can notice that the iterated QIP spline Kulkarni's type method improves over the QIP spline Kulkarni's type method.

Moreover, from Theorem 6 and Theorem 5, we can notice that the QIP spline Kulkarni's type method and the iterated QIP spline collocation method have the same order of convergence. However, from Proposition 7, we remark that the iterated QIP spline Kulkarni's type method improves over both iterated QIP spline collocation method and QIP spline Kulkarni's type method.

# 5. Numerical results

In this section, at first, we present results related to two integral equations of type (2.1), in order to give a numerical counterpart of the theoretical estimates given in the previous section. In fact, we have numerically solved the mentioned equations with QIP spline collocation method and QIP spline Kulkarni's type method in their basic and iterated versions, using both projectors  $Q_2$  and  $Q_3$ .

The integrals occurring in the various methods are evaluated by using the quadrature formulas of composite type presented in Appendix A, which are suitable in order to evaluate integrals with logarithmic kernel. For all the tests, for increasing values of n, we have computed the maximum absolute error by an approximate infinity norm calculated in this way

$$\|\varphi - \varphi_n\|_{\infty} := \max_{v \in G} |\varphi(v) - \varphi_n(v)|$$

where  $\varphi$  is the exact solution of the equation,  $\varphi_n$  is the approximate solution by one of the mentioned methods and G is a partition of the interval  $\mathcal{I}$  with mesh size h/7.

We have also computed the corresponding numerical convergence order, obtained applying the base 2 logarithm to the ratio between two consecutive errors.

Test 1

The first considered weakly singular Fredholm-Hammerstein integral equation reads

$$x(s) - \int_0^1 \log|t - s| x^2(t) dt = f(s), \qquad s \in \mathcal{I}$$

where

$$f(s) = -\frac{1}{9}\log(1-s) + \frac{1}{9}\log(1-s)s^9 - \frac{1}{9}\log(s)s^9 + \frac{1}{9}s^8 + \frac{1}{18}s^7 + \frac{1}{27}s^6 + \frac{1}{36}s^5 + \frac{46}{45}s^4 + \frac{1}{54}s^3 + \frac{1}{63}s^2 + \frac{1}{72}s + \frac{1}{81}s^7 + \frac{1}{27}s^6 + \frac{1}{36}s^5 + \frac{1}{45}s^6 + \frac{1}{54}s^3 + \frac{1}{63}s^5 + \frac{1}{63}s^5$$

We note that

$$\lim_{s \to 0} f(s) = \frac{1}{81}, \qquad \qquad \lim_{s \to 1} f(s) = \frac{29809}{22680}.$$

The exact solution of this equation is  $\varphi(s) = s^4$ .

By using Matlab environment, we have constructed the computational procedure in order to numerically solve this equation with the various methods presented in this paper and we have obtained the results reported in Table 1.

In the first, second and last column it is clear that the convergence order increases when the mesh size decreases. This fact is due to the term  $\log(h)$  that is present in the order of convergence of the corresponding methods, stated in Theorem 6, Proposition 7 and Theorem 5, whose effect in decreasing the convergence order is more prevailing when the mesh size h is wider.

As stated theoretically in the previous section, we underline that QIP spline Kulkarni's type method and iterated QIP spline collocation method are equivalent in terms of convergence order, but it is interesting to point out that the second one is easier to construct and it is cheaper in terms of computational cost. On the other side, the computational effort for the QIP spline Kulkarni's type method can be justified by the increased convergence order of the iterated QIP spline Kulkarni's type method. In fact, with a little additional computational effort with respect to its basic version, this last method can achieve a better order of convergence, as stated in Proposition 7 and as confirmed by the numerical results in the second column of Table 1.

Table 1: Numerical results for Test 1 with all presented methods and both projectors.

Test 1										
Kulkarni		Iterated Kulkarni		Collocation		Iterated collocation				
n	$\ \varphi-\varphi_n^k\ _{\infty}$	$O_{\infty}^{k}$	$\ \varphi - \tilde{\varphi}_n^k\ _{\infty}$	$\tilde{O}^k_\infty$	$\ \varphi-\varphi_n^c\ _{\infty}$	$O_{\infty}^{c}$	$\ \varphi - \tilde{\varphi}_n^c\ _{\infty}$	$\tilde{O}^c_{\infty}$		
	Methods based on $Q_2$									
2	3.68(-03)		2.08(-03)		3.26(-02)		7.18(-03)			
4	4.22(-04)	3.1	1.14(-04)	4.2	3.11(-03)	3.4	6.98(-04)	3.4		
8	3.47(-05)	3.6	4.93(-06)	4.5	3.68(-04)	3.1	5.82(-05)	3.6		
16	2.42(-06)	3.8	2.05(-07)	4.6	4.48(-05)	3.0	4.53(-06)	3.7		
32	1.58(-07)	3.9	8.19(-09)	4.6	5.51(-06)	3.0	3.37(-07)	3.8		
Methods based on $Q_3$										
4	1.38(-05)		4.71(-06)		2.41(-04)		6.57(-05)			
8	1.12(-06)	3.6	1.37(-07)	5.1	1.45(-05)	4.1	2.90(-06)	4.5		
16	5.03(-08)	4.5	4.38(-09)	5.0	8.90(-07)	4.0	1.11(-07)	4.7		
32	1.85(-09)	4.8	9.79(-11)	5.5	5.57(-08)	4.0	3.88(-09)	4.8		

#### Test 2

The second considered weakly singular Fredholm-Hammerstein integral equation reads

$$x(s) - \int_0^1 \log|s - t| \sqrt{1 + x^2(t)} dt = f(s), \qquad s \in \mathcal{I}$$

where

$$f(s) = \operatorname{Chi}(1-s)\sinh(s) + \sinh(s) + \sinh(s) + \sinh(s)\cosh(s) - \cosh(s) - \sinh(s) - h(s) -$$

$$Shi(s) = \int_0^s \frac{\sinh(t)}{t} dt, \qquad Chi(s) = \gamma + \log(s) + \int_0^s \frac{\cosh(t) - 1}{t} dt$$

and  $\gamma$  the Eulero-Mascheroni constant.

We point out that we consider f as a real function. We note also that

$$\lim_{s \to 0} f(s) \simeq 1.05725,$$
  $\lim_{s \to 1} f(s) \simeq 2.50031$ 

The exact solution of this equation is  $\varphi(s) = \sinh(s)$ .

In Table 2 we reported the results of numerical simulations related to this equation. Also these numerical tests confirm all the theoretical results and the remarks written previously.

Table 2: Numerical results for Test 2 with all presented methods and both projectors.										
		·		Test	2	·		·		
Kulkarni		Iterated Kulkarni		Collocation		Iterated collocation				
n	$\ \varphi - \varphi_n^k\ _{\infty}$	$O_{\infty}^{k}$	$\ \varphi - \tilde{\varphi}_n^k\ _{\infty}$	$\tilde{O}^k_\infty$	$\ \varphi - \varphi_n^c\ _{\infty}$	$O_{\infty}^{c}$	$\ \varphi - \tilde{\varphi}_n^c\ _{\infty}$	$\tilde{O}^c_{\infty}$		
	Methods based on $Q_2$									
2	2.47(-04)		6.50(-05)		1.50(-03)		2.86(-04)			
4	1.62(-05)	3.9	1.95(-06)	5.1	1.85(-04)	3.0	2.21(-05)	3.7		
8	1.04(-06)	4.0	6.20(-08)	5.0	2.29(-05)	3.0	1.69(-06)	3.7		
16	6.50(-08)	4.0	2.39(-09)	4.7	2.83(-06)	3.0	1.26(-07)	3.7		
32	4.04(-09)	4.0	8.48(-11)	4.8	3.50(-07)	3.0	9.10(-09)	3.8		
	Methods based on $Q_3$									
4	6.87(-07)		1.12(-07)		7.66(-06)		1.24(-06)			
8	3.14(-08)	4.5	3.04(-09)	5.2	5.84(-07)	3.7	5.29(-08)	4.6		
16	1.13(-09)	4.8	5.60(-11)	5.8	4.01(-08)	3.9	1.97(-09)	4.7		
32	3.80(-11)	49	9.21(-13)	5.9	2.62(-09)	3.9	6.85(-11)	4.8		

#### Test 3

Let us conclude this section, presenting some results related to a non smooth solution, as found in [10]. We consider the weakly singular Fredholm-Hammerstein integral equation

$$x(s) - \int_0^1 \log|t - s| x^2(t) dt = f(s), \qquad s \in \mathcal{I}$$

where f is chosen so that the exact solution is  $\varphi(s) = s \log(s)$ .

Numerical results are collected in Table 3: the presented approach shows an expected decay of its performance, if applied to an integral equation having a less regular solution. Moreover, the use of  $Q_3$  does not improve the order of convergence of  $Q_2$  as before, even if the errors are smaller. In any case, errors are in line with those presented in [10], where piecewise constant basis functions have been used on suitable graded meshes, instead of splines on classical uniform grids as employed here.

Even if not yet supported by the theory, these last simulations and analogous ones, not reported here, show the robustness of the proposed approach also in a non smooth framework.

Table 3: Numerical results for Test 3 with all presented methods and both projectors.									
Test 3									
Kulkarni		Iterated Kulkarni		Collocation		Iterated collocation			
$\frac{}{n \mid \ \varphi - \varphi_n^k\ _{\infty}}$	$O_{\infty}^{k}$	$\ \varphi - \tilde{\varphi}_n^k\ _{\infty}$	$\tilde{O}^k_\infty$	$\ \varphi-\varphi_n^c\ _{\infty}$	$O_{\infty}^{c}$	$\ \varphi - \tilde{\varphi}_n^c\ _{\infty}$	$\tilde{O}^c_{\infty}$		
Methods based on $Q_2$									
8   6.89(-04)		1.94(-04)		1.46(-02)		3.87(-03)			
16 $9.37(-05)$	2.9	9.08(-06)	4.4	6.32(-03)	1.2	5.25(-04)	2.9		
$32 \mid 1.37(-05)$	2.8	4.22(-07)	4.4	3.29(-03)	0.9	7.76(-05)	2.8		
Methods based on $Q_3$									
$8 \mid 6.23(-05)$		6.48(-06)		1.00(-02)		4.30(-04)			
$16 \mid 1.01(-05)$	2.6	4.29(-07)	3.9	4.92(-03)	1.0	5.82(-05)	2.9		

4.3

2.44(-03)

1.0

8.68(-06)

2.7

Table 3: Numerical results for Test 3 with all presented methods and both projectors

### 6. Conclusions

In this paper, spline quasi-interpolating projectors have been used to efficiently solve nonlinear Fredholm-Hammerstein integral equations with logarithmic kernel by means of collocation and Kulkarni methods, both in their basic and iterated versions. Theoretical analysis of discretization error and convergence order has been provided, and numerical results have been shown validating the estimates, obtained under the hypothesis of sufficiently smooth solutions. The analysis of the proposed approach performance in a non-smooth framework is currently under study, but related numerical results appear promising.

Moreover, the methodologies proposed in this paper can be extended to an integral operator of the form

$$K(x)(s) := \int_0^1 r(s,t) \log |s-t| \, \psi(t,x(t)) dt, \qquad s \in \mathcal{I}, \qquad x \in X,$$

where r(s,t) is a smooth function defined in  $\mathcal{I} \times \mathcal{I}$ .

Future investigations will be devoted to treat kernels with higher order of singularity, such as those giving rise to Cauchy principal value or Hadamard finite part integrals (the reader is referred, for instance, to [2, 3] for some examples of such types of kernels arising in BEMs).

# Appendix A. Quadrature formulas

1.61(-06)

Throughout this work we met several times the following type of integrals

2.22(-08)

$$\int_{a}^{b} \log|t - s| f(t)dt, \quad s \in [a, b].$$

By reader's convenience, in this paragraph we briefly recall suitable quadrature formulas in order to numerically calculate them in an efficient way (see [2, 3]).

We start by considering a = -1, b = -1 and the application of the following interpolating quadrature formula

$$\int_{-1}^{1} \log|t - s| f(t) dt \approx \sum_{k=1}^{n} \omega_k \tilde{\gamma}_k(s) f(x_k),$$

where  $x_k$  and  $\omega_k$  are respectively the knots and the weights of the classical n points Gauss-Legendre quadrature formula in [-1,1] and  $\tilde{\gamma}_k(s)$  are defined by

$$\tilde{\gamma}_k(s) = \frac{1}{2} \sum_{i=0}^{n-1} (2i+1)\mu_i(s) P_i(x_k).$$

In this equality  $P_i(t)$  is the Legendre polynomial of degree i and  $\mu_i(s)$  are the modified moments of the kernel  $\log |t-s|$ , defined as

$$\mu_i(s) = \int_{-1}^1 \log|t - s| P_i(t) dt.$$

They can be computed by the following recursive procedure:

$$\begin{cases} \mu_0(s) = (1+s)\log(1+s) + (1-s)\log(1-s) - 2\\ \mu_j(s) = \frac{1}{2j}\Theta_{j-1}(s), \quad j \ge 1 \end{cases}$$

where:

$$\begin{cases} \Theta_0(s) = (1 - s^2) \log \left(\frac{1 - s}{1 + s}\right) - 2s \\ \Theta_1(s) = 2s\Theta_0(s) + \frac{8}{3} \\ \Theta_j(s) = \frac{j + 1}{j(j + 2)} [(2j + 1)s\Theta_{j-1}(s) - j\Theta_{j-2}(s)], \qquad j \ge 2 \end{cases}$$

Regarding the degree of accuracy of these quadrature formulas, the reader can refer to [3].

Using standard techniques, such as change of variable or subdivision of the integration domain, we can construct formulas in order to calculate this type of integrals over an interval [a,b], and also composite formulas.

For sake of completeness we briefly construct the composite quadrature formula for this type of integrals over [a, b]. We point out that this formula has been used throughout this work, all times where an integral of such type has occurred.

Defining a positive integer m and setting h = (b - a)/m, we set a uniform partition of the interval [a, b], made by m subintervals

$$Z = \{z_{\eta} = a + \eta h, \eta = 0, \dots, m\}.$$

So we have

$$\int_{a}^{b} \log|t - s| f(t) dt = \sum_{n=0}^{m-1} \int_{z_n}^{z_{n+1}} \log|t - s| f(t) dt.$$

Fixing s, we call  $\tilde{\eta}$  the index such that  $s \in [z_{\tilde{\eta}}, z_{\tilde{\eta}+1}]$ . We note that the only singular integral is the one on the domain  $[z_{\tilde{\eta}}, z_{\tilde{\eta}+1}]$ . So only this integral must be calculated using the formula outlined in this section, while for the other subintervals we can use a classic Gauss-Legendre formula. After a change of variable and some algebra we reach

$$\int_{a}^{b} \log|t-s|f(t)dt \approx \frac{h}{2} \left[ \log\left(\frac{h}{2}\right) \sum_{k=1}^{n} \omega_{k} f(x_{k,\tilde{\eta}}) + \sum_{k=1}^{n} \omega_{k} \tilde{\gamma}_{k}(\sigma) f(x_{k,\tilde{\eta}}) + \sum_{\eta=0, \eta\neq\tilde{\eta}}^{m-1} \sum_{k=1}^{n} \omega_{k} f(x_{k,\eta}) \log|x_{k,\eta}-s| \right],$$

where 
$$x_{k,*} = \frac{h}{2}x_k + \frac{z_* + z_{*+1}}{2}$$
 with  $* = \eta, \tilde{\eta}, \sigma = -\frac{2}{h} \left( \frac{z_{\tilde{\eta}} + z_{\tilde{\eta}+1}}{2} - s \right)$ .

# Appendix B. Proof of Theorem 1 for $Q_3$

In this appendix we prove that the QIP  $Q_3$  satisfies Theorem 1 and therefore Theorem 2 holds.

Let  $\mathbb{P}_d$  be the space of polynomials of degree at most d. Consider the interval  $[t_{i-1}, t_i]$ ,  $i = 4, \ldots, n-4$ , the middle point  $s_i$ , defined as in Section 3 and  $m_4(t) = t^4$ . Therefore, we can write  $m_4(t) = (t - s_i)^4 + p_3(t) = p_4(t) + p_3(t)$ , where  $p_3 \in \mathbb{P}_3$ . As  $Q_3p_3 = p_3$ , we can write

$$\int_{t_{i-1}}^{t_i} (Q_3 m_4(t) - m_4(t)) dt = \int_{t_{i-1}}^{t_i} (Q_3 p_4(t) - p_4(t)) dt.$$

Now, as  $\int_{t_{i-1}}^{t_i} p_4(t)dt = \frac{h^5}{80}$ , it is sufficient to prove that also  $\int_{t_{i-1}}^{t_i} Q_3 p_4(t)dt$  is equal to  $\frac{h^5}{80}$ .

From the expression of the coefficient functionals  $\lambda_i(x)$ ,  $i = d+1, \ldots, n$  of  $Q_3$  given in [19], it is possible to obtain the quasi-Lagrange form (3.3) of  $Q_3$ . Therefore

$$\int_{t_{i-1}}^{t_i} Q_3 p_4(t) dt = \int_{t_{i-1}}^{t_i} \sum_{j=0}^{2n} (\xi_j - s_i)^4 L_j(t) dt = \sum_{j=0}^{2n} (\xi_j - s_j)^4 \int_{t_{i-1}}^{t_i} L_j(t) dt.$$

Taking into account the locality of the B-splines, the symmetry of the data points with respect to  $s_i$  and the symmetry properties of the coefficients  $\lambda_i(x)$ ,  $i=d+1,\ldots,n$ , we can compute  $\int_{t_{i-1}}^{t_i} L_j(t)dt$  and after some algebra we deduce  $\int_{t_{i-1}}^{t_i} Q_3 p_4(t)(t) = \frac{h^5}{80}$ . Therefore, considering the QIP  $Q_3$ , Theorem 1 holds also for the odd case d=3.

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# References

- [1] M. Abramowitz, I.A. Stegun, Handbook of Mathematical Functions, Dover Publications, New York, 1964.
- [2] A. Aimi, M. Diligenti, G. Monegato, New numerical integration schemes for applications of Galerkin BEM to 2D problems, Int. J. Numer. Methods Eng., 40.11 (1997), 1977–1999.
- [3] A. Aimi, M. Diligenti, G. Monegato, Numerical integration schemes for the BEM solution of hypersingular integral equations, Int. J. Numer. Methods Eng., 45.12 (1999), 1807–1830.
- [4] A. Aimi, M. Diligenti, M.L. Sampoli, A. Sestini, Isogemetric analysis and symmetric Galerkin BEM: a 2D numerical study, *Applied Mathematics and Computation*, **272** (2016), 173–186.
- [5] C. Allouch, S. Remogna, D. Sbibih, M. Tahrichi, Superconvergent methods based on quasi-interpolating operators for fredholm integral equations of the second kind, Appl. Math. Comput., 404 (2021), 1-14.
- [6] C. Allouch, P. Sablonnière, D. Sbibih, Solving Fredholm integral equations by approximating kernels by spline quasiinterpolants. *Numer. Algorithms*, 56 (2011), 437–453.
- [7] C. Allouch, P. Sablonnière, D. Sbibih, A modified Kulkarni's method based on a discrete spline quasi-interpolant. Math. Comput. Simul., 81 (2011), 1991–2000.
- [8] C. Allouch, P. Sablonnière, D. Sbibih, A collocation method for the numerical solution of a two dimensional integral equation using a quadratic spline quasi-interpolant. *Numer. Algorithms*, 62 (2013), 445–468.
- [9] C. Allouch, P. Sablonnière, D. Sbibih, M. Tahrichi, Product integration methods based on discrete spline quasi-interpolants and application to weakly singular integral equations, J. Comp. Appl. Math., 233, (2010), 2855-2866.
- [10] C. Allouch, D. Sbibih, M. Tahrichi, Numerical solutions of weakly singular Hammerstein integral equations, Appl. Math. Comput., 329 (2018), 118 – 128.
- [11] P. Assari, Thin plate spline Galerkin scheme for numerically solving nonlinear weakly singular Fredholm integral equations, Applicable Analysis, (2018), https://doi.org/10.1080/00036811.2018.1448073.
- [12] P. Assari, M. Dehghan, A meshless method for the numerical solution of nonlinear weakly singular integral equations using radial basis functions, Eur. Phys. J. Plus (2017), 132: 199.

- [13] K.E. Atkinson, A survey of numerical methods for solving nonlinear integral equations, J. Integral Eqns Appl., 4 (1992), 15 46.
- [14] K.E. Atkinson, G. Chandler, Boundary integral equation methods for solving Laplace's equation with nonlinear boundary conditions: the smooth boundary case, *Math. of Comp.*, **55**(192), (1990), 451-472.
- [15] K.E. Atkinson, W. Han, Theoretical Numerical Analysis, A Functional Analysis Framework (Third Edition), Springer, New York, 2009.
- [16] D. Barrera, F. Elmokhtari, D. Sbibih, Two methods based on bivariate spline quasi-interpolants for solving Fredholm integral equations, *Applied Numerical Mathematics*, **127** (2018), 78 94.
- [17] C. de Boor, A Practical Guide to Splines, Revised Edition, Springer, New York, 2001.
- [18] A. Carabineanu, A boundary element approach to the 2D potential flow problem around airfoils with cusped trailing edge, Comput. Methods Appl. Mech. Eng., 129(3) (1996), 213–219.
- [19] C. Dagnino, A. Dallefrate, S. Remogna, Spline quasi-interpolating projectors for the solution of nonlinear integral equations, J. Comput. Appl. Math., 354 (2019), 360 372.
- [20] C. Dagnino, S. Remogna, Quasi-interpolation based on the ZP-element for the numerical solution of integral equations on surfaces in  $\mathbb{R}^3$ . BIT Numer. Math., 57 (2017), 329–350.
- [21] C. Dagnino, S. Remogna, P. Sablonnière, On the solution of Fredholm integral equations based on spline quasi-interpolating projectors, *BIT Numer. Math.*, **54** (2014), 979 1008.
- [22] Z. Gouyandeh, T. Allahviranloo, A. Armand, Numerical solution of nonlinear Volterra-Fredholm-Hammerstein integral equations via Tau-collocation method with convergence analysis, *J. Comput. Appl. Math.*, **308** (2016), 435–446.
- [23] L. Grammont, R.P. Kulkarni, T.J. Nidhin, Modified projection method for Urysohn integral equations with non-smooth kernels, J. Comput. Appl. Math., 294 (2016), 309–322.
- [24] L. Grammont, R.P. Kulkarni, P.B. Vasconcelos, Fast and accurate solvers for weakly singular integral equations, *Numer. Algorithms* https://doi.org/10.1007/s11075-022-01376-x.
- [25] H. Kaneko, R.D. Noren, Y. Xu, Numerical solutions for weakly singular Hammerstein equation and their superconvergence, J. Int. Eq. Appl., 4(3) (1992), 391–407.
- [26] H. Kaneko, R.D. Noren, P.A. Padilla, Superconvergence of the iterated collocation methods for Hammerstein equations, J. Comput. Appl. Math., 80 (1997), 335–349.
- [27] H. Kaneko, Y. Xu, Superconvergence of the iterated Galerkin methods for Hammerstein equations, SIAM J. Numer. Anal., 33 (1996), 1048–1064.
- [28] R.P. Kulkarni, A superconvergence result for solutions of compact operator equations. B. Aust. Math. Soc., 68(3) (2003), 517–528.
- [29] D. Lesnic, L. Elliott, D.B. Ingham, Boundary element methods for determining the fluid velocity in potential flow, Eng. Anal. Boundary Elem., 11(3) (1993), 203–213.
- [30] M. Mandal, G. Nelakanti, Superconvergence Results for Weakly Singular Fredholm-Hammerstein Integral Equations, Numer. Funct. Anal. Optim., 40 (2019), 548-570.
- [31] G. Monegato, V. Colombo, Product integration for the linear transport equation in slab geometry, *Numer. Math.*, **52** (1988), 219–240.
- [32] A. Pedas, G. Vainikko, The Smoothness of Solutions to Nonlinear Weakly Singular Integral Equations, *Journal for Analysis and its Application*, **13**(3) (1994), 463–476.
- [33] B.T. Polyak, Newton-Kantorovich method and its global convergence, J. Math. Sci., 133, 1513-1523 (2006).
- [34] L.L. Schumaker, Spline functions basic theory. Cambridge Mathematical Library, 1981.
- [35] A.M. Wazwaz, Linear and Nonlinear Integral Equations, Methods and Applications, Higher Education Press, Beijing and Springer-Verlag Berlin Heidelberg, 2011.